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# ON THE BEHAVIOUR OF THE IONOSPHERE DURING SUDDEN IONOSPHERIC DISTURBANCES.

The ionization equation was solved for the electron production function linear with time. It is indicated that the development of a flare in  $H_\alpha$  light is not representative for an active radiation of flares. The method of determination of the value  $\tau_0 = \frac{1}{\frac{dN}{dt} - \frac{N}{N_0}}$  and of the course with time of solar flare radiation active for the ionosphere is suggested.

## I.

Relative electron density variations during sudden ionospheric disturbances (SID'S) due to active emission of solar flares are described by the well-known equation:

$$\tau_0 = \frac{d}{dt} \frac{N(t)}{N_0} = \frac{I(t)}{I_0} - \left( \frac{N(t)}{N_0} \right) \quad (1)$$

Here  $N_0$  is electron density and  $I_0$  - function of the electron production (both - at the moment of SID beginning).

Recently a solution of equation (1) has been found [1,2] for  $I(t)/I_0$  in the form of a rectangular pulse. If is a linear function of time, i.e.  $I(t)/I_0 = P_0 + \frac{t}{T}$

$$\frac{N(t)}{N_0} = \frac{1}{\beta^{1/3}} \cdot \frac{[u^{10} - \alpha \beta^{1/3} \cdot v^0] \cdot [u - u^{10} - \alpha \beta^{1/3} u^0] \cdot v'}{v^{10} - \alpha \beta^{1/3} \cdot v^0 \cdot [u - u^{10} - \alpha \beta^{1/3} u^0] \cdot v} \quad (2)$$

where  $\beta = \frac{T}{\tau_0}$ ;  $T$  is a parameter characterizing the variations of  $I(t)/I_0$ ; the values of  $T > 0$  correspond to an increasing  $I(t)/I_0$  and  $T < 0$  - to a decreasing one as  $\frac{N(0)}{N_0}$  for intervals having different values of  $T$ ;  $u, v, u', v'$  are the Airy functions and their derivatives with argument  $Z = \beta^{2/3} (P_0 + \frac{t}{T})$  and  $u^0, v^0, u^{10}, v^{10}$  are the same ones with argument  $Z^0 = \beta^{2/3} P_0$

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Lets us consider some applications of these relations.

It seems interesting to find out if the development of a flare in  $H_{\alpha}$  light is representative for an active flare emission /4,5/. The development curves of flares (the course with time of intensity and emission flow in  $H_{\alpha}$  light) were approximated by the broken lines.

Electron density variations were computed by means of equation (2) assuming that time variations of active emission are the same as in  $H_{\alpha}$  light, but with a greater amplitude. These calculations were made for flares 31/VIII-56, 28/VIII-57, 3/IX-57 and 23/XI-57. On the other hand, electron density variations were determined by known measured values of  $f_{min}$  /6/ to compare them with observations. Fig.1 gives an example of comparison of observed values of  $N(t)/N_0$  with those computed by a described method. Graphs of fig. 1 shows that the development of the flare in  $H_{\alpha}$  light is not representative for the active emission. Apparently this emission is X-ray radiation of the flare. The reverse problem is also of a great interest. The ionospheric parameter  $\tau_0$  and the course with time of active emission can be determined if there are continuous measurements of  $N(t)/N_0$  during SIE (for D- region by measurements of absorption,  $f_{min}$  e.t;c. ) Really, supposing that an active emission in initial stage of development of the flare can be approximated by triangular pulse we can find  $\tau_0$  and parameters determining that triangle, they are:  $t_m$  - time delay of electron density maximum  $N_m/N_0$  relatively to the top of the triangle;  $T_0$  and  $T_1$  - values of  $T$  for different sides of triangle;  $t_0$  - the total time of  $I(t)/I_0$  growth. To solve this problem it is necessary to have five equations:

$$\frac{N_m}{N_0} = \Psi(T_0, t_0, T_1, \tau_0) \quad (3a)$$

$$\frac{N_m}{N_0} = \sqrt{1 + \frac{t_0}{T_0} + \frac{t_r}{T_1}} \quad (3b)$$

$$t_r = t_0 + t_r \quad (3c)$$

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$$\frac{d^2}{dt^2} \left( \frac{N}{N_0} \right) \Big|_{t=0} = \frac{1}{T_0^2} \quad (3d)$$

$$\frac{d^2}{dt^2} \left( \frac{N}{N_0} \right) \Big|_{t=t_m} = \frac{1}{T_0 T_1} \quad (3e)$$

where  $T_0 \neq 0$  and  $T_1 \neq 0$ . The equation (3a) is an equation (2) for the moment  $t_m$ . The relation (3b) is found from the condition  $\frac{d}{dt} \left( \frac{N}{N_0} \right) \Big|_{t=t_m} = 0$ . In (3c)  $\theta$  is

the time, measured from the beginning of SID to the moment of  $\frac{N(t)}{N_0}$  maximum. The equations (3d) and (3d) and (3e) for the moments  $t=0$  and  $t=t_m$  respectively are obtained by calculations of the second derivatives of  $N(t)/N_0$ ; their meanings must be taken from experimental graph  $N(t)/N_0$ . The value  $T_0$  can be found from solution of equations (3); the course with time  $I(t)/I_0$  can be determined by this method of calculation. It must be noted that the course with time of  $N(t)/N_0$  during SID must be continuous and sufficiently smooth.

The course with time of  $I(t)/I_0$  for the flare 28/VIII-57, is presented on fig.2 calculated by just described method. Variations of  $N(t)/N_0$  during SID were determined by  $f_{min}$ ; the smoothed values of  $f_{min}$  were used. It was found, that

$$T_0 = 4800 \text{ sec. and } T_1 = 600 \text{ sec.}$$

Thus, taking into account the nonstationary process during SID  $1/s$  the method of determination for value  $T_0 = \frac{1}{\alpha N_0}$  and course with time of solar flares emission, active for ionosphere is suggested.

It seems possible to apply this method to the determination of an effective recombination coefficient using the diurnal changes of critical frequencies for different ionospheric layers.

It would be also very interesting to compare such

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analysis of SID with the data on X-ray radiation of flares, obtained on the artificial satellites of the Earth.

#### R E F E R E N C E S

1. J. Taubenheim, J. Atm. Terr. Phys., vol. 11, No. 1, 1957, 14-22.
2. Н.А. Савич Известия Крымской Астрофизической Обсерватории. Т. 19.
3. В.А. Фок. Таблицы функции Эйри, 1946 г. Москва.
4. M. A. Ellison, Solar eclipses and the ionosphere. 1956, 180-183.
5. V. C. A. Ferraro, Nature, vol. 175, No. 4449, 1955, 242-244.
6. E. V. Appleton, J. Atm. Terr. Phys., Vol. 3, No. 5, 1953, 282-284.

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CAPTIONS.

Fig. 1 Comparison of the experimental and theoretical curves of the electron density variations during the flares.

Time - Moscow;  $\tau_0 = 4000$  sec.

a) 1. Intensity variations  $I(t)/I_0$  of the flare in H

2. Variations  $\frac{N(t)}{N_0}$  exp determined by  $f_{min}$

3. Theoretical variations  $\frac{N(t)}{N_0} / th$

b,c) 1. Variations of the emission flow  $\phi(t)/\phi_0$  of the flare in H $\alpha$

2. Variations  $\frac{N(t)}{N_0} / exp$  determined by  $f_{min}$ .

3. Theoretical variations  $\frac{N(t)}{N_0} / th$

Fig.2 Variations of electron production function  $I(t)/I_0$  during the flare 28/VIII-57

1. Variations  $\frac{N(t)}{N_0} / exp$ - determined by  $f_{min}$

2. Smoothed curve  $\frac{N(t)}{N_0} / exp.$

3. Approximating triangular function  $\frac{I(t)}{I_0}$

4. Variations  $I(t)/I_0$  computed by eq. (1) with  $\tau_0 = 4800$  sec.